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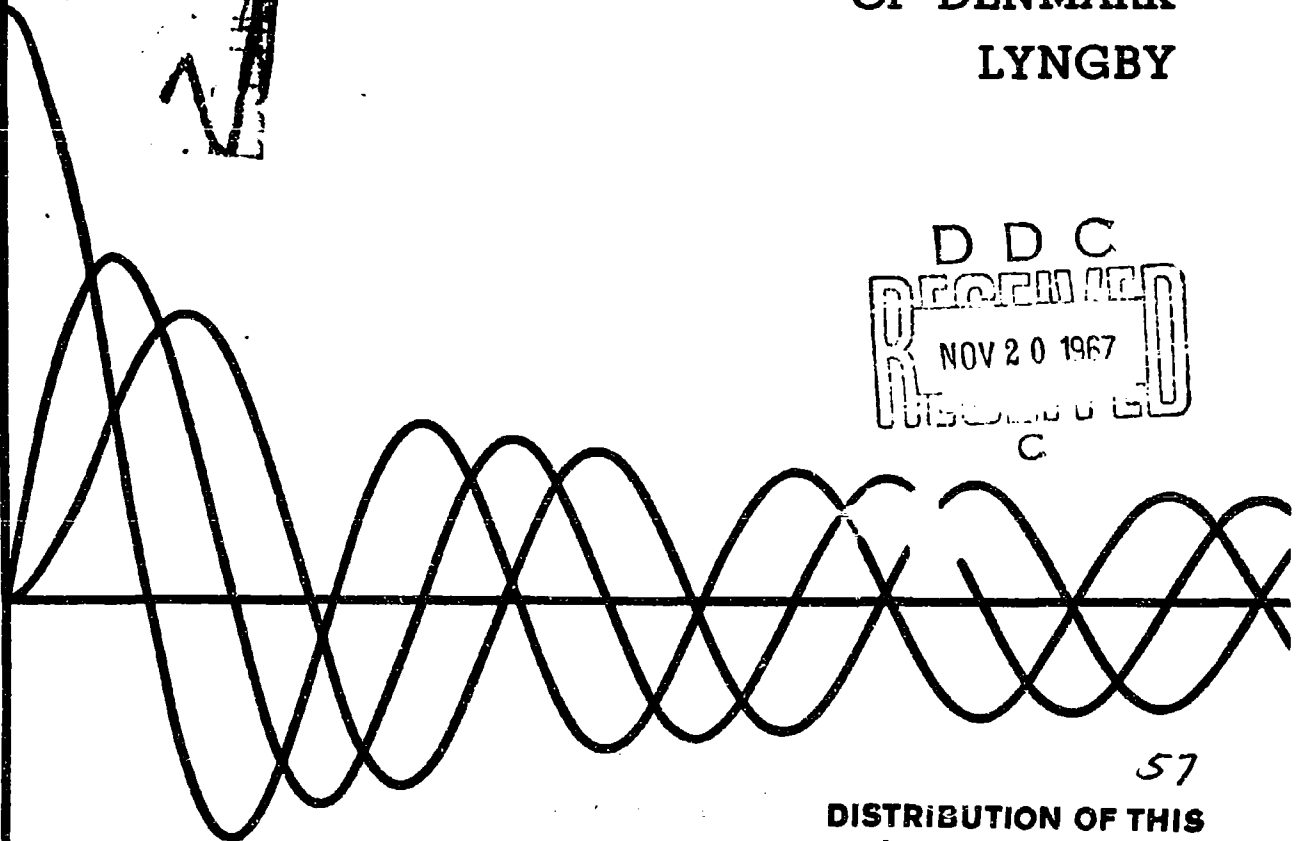
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Contract AF 61(052)-794

FINAL SCIENTIFIC REPORT

"Reflector Arrays"

1 April 1964 - 30 June 1967

E. Dragg Nielsen

July 1st, 1967

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ABSTRACT

The Van Atta reflector was first described in a patent by Dr. L.C. Van Atta in 1959. The advantage of this passive reflector type should be that the reradiated field has a maximum back in the direction of arrival of the primary plane wave. Since this retrodirective effect of the reflector might be of great importance if used as a navigational aid in the air or at sea, it seemed worth while to carry out a theoretical investigation of such reflectors, especially since only experimental investigations had been made before this contract was initiated.

The work performed under the contract deals mainly with theoretical and numerical investigations of Van Atta reflectors consisting of dipoles. A survey of the literature concerning active or passive Van Atta reflectors has been made. Both a linear and a two-dimensional plane Van Atta reflector has been investigated numerically and a theory for arbitrary Van Atta reflectors has been developed. An experimental investigation of a linear Van Atta reflector was carried out and the results compared with the theoretical results.

TABLE OF CONTENTS

Abstract	1
1. Introduction	3
2. Survey of literature	5
3. Theoretical investigation of arbitrary Van Atta arrays	6
4. A linear Van Atta reflector consisting of four half-wave dipoles	8
5. Square Van Atta reflector with or without a conducting plate	11
6. Bandwidth properties of the square Van Atta reflector	13
7. Conclusion	15
8. Computer program	16
9. References	17
Appendix 1	20
Figures	

1. INTRODUCTION

The contract AF 61(052)-794 was running from April 1, 1964 to June 30, 1967.

The objectives of the contract were given as follows:

Carry out a theoretical investigation of the performance of a Van Atta reflector and, by electronic computer calculations, to find the radiation patterns of a number of such reflectors with various characteristics. It is expected that this would lead to a procedure for the design of a Van Atta reflector with a prescribed radiation pattern. It is the intention first to treat the linear and two-dimensional array, later the method of investigation may be extended to include Van Atta reflectors placed on circular cylinders and spheres. Further efforts will be made to supplement the theoretical investigation by experiment.

In this final report a summary of all the work performed under the contract is given. Some related work carried out at this laboratory but not paid by the contract is described, too.

First an analytical investigation of arbitrary Van Atta reflectors was carried out. Each pair of antenna elements with connecting transmission lines was described by an equivalent circuit. Reradiation from the elements and mutual interaction were taken into account. Next a theoretical investigation of a linear Van Atta reflector consisting of four parallel half-wave dipoles was performed. The dipole elements were equispaced and mutual impedances between the dipoles were taken into account. An expression for the reflected field was set up by superposition of the three fields involved. Further an analysis of the general shape of the reradiation pattern was made neglecting the mutual impedances. This analysis showed that for some combinations of the parameters involved the reflector does not act as stated in the patent description of Van Atta reflectors.

An experimental investigation of the four element linear Van Atta reflector was carried out in a radio-anechoic chamber. In agreement with the theoretical investigation, the experiments showed that the reflector only to a limited extent has the retrodirective effect stated in the patent description and that it has a mirror effect to the same extent as it has a retrodirective effect.

Numerical investigations of a four element Van Atta reflector and an optimization of the reflector with respect to the parameters involved was carried out using an electronic computer. The reradiation pattern of the reflector was optimized with the criterium that the minimum values of the

back-scattered field intensities for all angles of incidence should be as large as possible.

Finally a square Van Atta reflector consisting of half-wave dipoles was investigated theoretically and numerically. The effect of mounting the dipoles above and parallel to a conducting plate was also examined. The mutual impedance between the dipole elements and the reradiation both from the elements and from the plate was taken into account. All parameters involved have been varied in a numerical analysis of the reflector. The parameters are: the number of dipole elements, the length of the dipoles, the length and characteristic impedance of the transmission lines, the distance between the dipoles, and the distance from the dipoles to the conducting plate. The effect of changing the parameters is shown in curves of the back-scattering cross section of the reflector as a function of the variation of each of the parameters. The numerical results have been compared with experimental results obtained by others.

The last work on the contract has been an investigation of the bandwidth properties of a 4 by 4 element square Van Atta reflector consisting of dipoles.

2. SURVEY OF LITERATURE

Since Van Atta ¹⁾ in 1959 proposed his passive retrodirective reflector several papers have suggested the use of this reflector type in various communication systems ^{2) - 23)}. Some of the papers tend to give an analytical treatment of the Van Atta reflector but most of them neglect the scattering by and coupling between the antennas. Many of the papers suggest active components inserted in the transmission lines of the reflector, such as modulated phase shifters ⁵⁾, amplifiers ^{6) 10)}, and mechanical modulation by means of cavity resonators ⁹⁾. It has been suggested to use Van Atta reflectors for satellite communication and both passive, semipassive and active systems have been proposed ^{6) 15) 21) 22)}. Other applications are for navigational aids, for example used to enhance the reflection from radar targets, from lifeboats, and from aircraft ^{3) 5) 7) 10)}. A more detailed discussion of the literature concerning Van Atta reflectors and their applications is given in Scientific Report No. 1 and by Appel-Hansen ^{34) 35)}.

3. THEORETICAL INVESTIGATION OF ARBITRARY VAN ATTA ARRAYS

A general treatment which may be used for a number of different investigations is described in Scientific Report No. 1 of the contract. The configuration investigated is shown in fig. 1. The elements are supposed to be dipoles, but the theory could easily be extended to other antenna types. The dipoles are placed on and parallel to an imaginary smooth surface which may be f.ex. a plane, a cylinder or a sphere. The field incident on the reflector is a plane wave.

The open circuit voltage induced at the terminals of each antenna element by the primary plane wave is calculated. Using these voltages a system of linear equations is developed for calculating the currents in each antenna taking into account the mutual impedances, the characteristics of the interconnecting transmission lines, and the induced voltage at the element itself (giving the scattered field) and at its mate (giving the retrodirective field).

When the currents are determined the reradiated field may be calculated. For a reflector with all elements parallel this is done using the theory of antenna arrays. Finally the properties of the reflector array are described by calculation of the differential scattering cross section.

The mutual impedances of the dipoles are calculated using the induced EMF method as given in Jordan's book ³⁸⁾. Algol procedures for computing the mutual impedances between linear dipoles with sinusoidal current distributions and for arbitrary wire-antennas with a known current distribution has been developed at this laboratory.

The transmission lines are represented by equivalent circuits of the general x-circuit type. This type of equivalent circuit has been chosen because it has the advantage of being valid for all lengths of the transmission lines. It is assumed to be symmetrical and lossless.

Since both the induced voltages and the mutual impedances are complex quantities the real matrix equation to be solved will be of the $2N$ th order where N is the total number of antenna elements. This means that it will be almost necessary to use an electronic computer for solving the matrix equation numerically when reflectors with more than two elements are treated.

A simple numerical example of a four element Van Atta array with equispaced half-wave dipoles is given in SR 1 in order to illustrate the theory developed. The numerical results indicate that a retrodirective effect as

stated by Van Atta is obtained to some degree. However the results also show that the mirror effect of the reflector is of the same order of magnitude as the retrodirective effect. The influence of the mutual impedances and of a mismatch between the antenna elements and the transmission lines may be utilized to change the reradiation pattern to compare better with a prescribed form.

4. A LINEAR VAN ATTA REFLECTOR CONSISTING OF FOUR HALF-WAVE DIPOLES

A detailed investigation of a four element linear Van Atta reflector consisting of half-wave dipoles was carried out and is described in Scientific Report Nos. 2, 3, and 4 of the contract.

First an expression for the reflected field is developed considering each dipole as a receiving and transmitting antenna matched to its transmission line. When a plane wave is incident on the reflector a current distribution consisting of three different parts will be generated in the antenna elements. The first part is due to the energy transmitted to the antenna through the transmission line from its mate. The second part is the current induced in the antenna by the incident plane wave and the third part is induced in the antenna due to its mutual interaction with the other antennas of the reflector. Only the first part of the current distribution is considered in Van Atta's patent description, while the last two parts are neglected.

In fig. 2 is shown the retrodirective effect of the reflector due to the first part of the current distribution mentioned above. The second part of the current distribution creates a mirror effect of the reflector which when the dipoles are matched to the transmission lines is of the same order of magnitude as the retrodirective effect. This mirror effect which is shown in fig. 3 is not mentioned in the patent description of the reflector.

It is shown that the length of the transmission lines and the angle of incidence has a great influence on the reradiation, and for some values of these two parameters the first and second part of the current in the antennas are of opposite phase and they cancel each other so that the only reflected energy is the small part due to the mutual interaction between the elements. The length of the transmission lines at which the behaviour of the reflector is as much as possible in accordance with the patent description has been found but even this reflector has the mirror effect mentioned above.

Numerical analysis have been made with a spacing between the dipole elements of half a wavelength. It turns out that in most cases maximum reflection is not back in the direction of incidence of the primary plane wave, and that the mutual interaction between the dipoles causes asymmetries in the reradiation pattern. This means that it is possible to increase the back-scattered energy by choosing a proper combination of the distance between the elements and the length of the transmission lines. The above

mentioned investigations are described in Scientific Report No. 2.

In order to verify the theoretical results obtained, an experimental investigation of a four element linear Van Atta reflector was carried out in a radioanechoic box at the laboratory. This investigation is treated in Scientific Report No. 3. The four half-wave dipoles of the experimental reflector were slot fed dipoles with open-ended terminations. The length of the connecting transmission lines could be changed by means of line-stretchers thus examining the influence of the length of the lines on the reflecting properties of the reflector. The measurements were performed at 3.21 GHz as this frequency gives the best matching between each dipole and the connected transmission line.

The Van Atta reflector was placed on a moveable pedestal in the center of the anechoic box, and the measurements were based upon the principle of interference between the signal reflected by the reflector and a reference signal. Radiation patterns were measured in the plane normal to the axis of the dipoles for discrete values of the angle of incidence of the primary plane wave.

A good agreement between the experimental and theoretical results was found. The measured results confirmed the theoretical results of Scientific Report No. 2 and showed (1) that maximum reradiation is not always back in the direction of incidence, (2) that the reflector has a mirror effect to the same degree as it has a retrodirective effect, (3) that the reradiation depends strongly upon the length of the transmission lines, and (4) that the mutual impedances causes asymmetries in the radiation patterns.

As an extension of the theoretical analysis of the four element linear Van Atta reflector an optimization of this reflector has been carried out, the results of which is given in Scientific Report No. 4.

First the original expression for the reradiated field as derived in Scientific Report No. 2 was changed to a more general form which makes it possible to study the influence of asymmetries in the location of the dipoles, unequal lengths of the transmission lines, and a mismatch between the dipoles and the transmission lines. Further, a method was developed for computing a quantity which may be used as a measure of the deviation from the retrodirective effect of the reflector. This quantity is shown to be just as useful as the reradiation pattern itself when two different Van Atta reflectors are to be compared.

A perfectly working retrodirective Van Atta reflector has not been found. However, when mutual coupling is neglected, a condition for the smallest deviation from retrodirectivity has been derived. It is shown that the mutual coupling between the dipoles usually causes the reradiation of

the reflector to decrease and the deviation from retrodirectivity to increase. However for certain values of the parameters involved it turns out that coupling may increase the back-scattering up to 50% for some angles of incidence.

The numerical optimization of the reradiation pattern of the four element Van Atta reflector was carried out using a computational technique starting with an a priori reasonable set of parameters selected by examining 1600 different reflectors. The parameters were then perturbed about their initial values, the effect on the reradiation pattern was observed, and a new set of parameters giving an improved result was selected. The success of this method depends on the correctness of the original set of parameters and the computer program for perturbing the parameter values

An attempt was made to fulfil the following two criteria:

the minimum value of the back-scattered field intensity, as a function of the angle of incidence, should be as large as possible and the deviation from Van Atta effect as small as possible,
the minimum value of the back-scattered field intensity, as a function of the angle of incidence, should be above various prescribed levels and the deviation from Van Atta effect as small as possible.

For both optimization processes it turned out that the optimum value of the spacing was close to 1.5 wavelengths. Further it was found that, due to coupling, the minimum value of the back-scattered field intensity may be increased and the deviation from the retrodirective effect decreased if the transmission lines are permitted to be of unequal lengths and asymmetries are permitted in the location of the dipoles around the center of the reflector. However, the improvements are small and asymmetries often causes the opposite effect.

5. SQUARE VAN ATTA REFLECTORS WITH OR WITHOUT A CONDUCTING PLATE

Another configuration of the reflector array is the plane, square Van Atta reflector consisting of parallel dipoles. The investigation of this reflector is described in Scientific Report No. 5.

The theoretical investigation of this reflector has already been described in Scientific Report No. 1. However, in order to be able to compare the theory with experimental results obtained by Sharp²⁾ the effect of mounting the dipoles above and parallel to a conducting plate has to be taken into account. The system investigated is shown in fig. 4.

The reflecting properties of the plate are supposed not to be influenced by the presence of the dipoles. The reflected field is found using the method of physical optics as described f.ex. in Kerr's book³⁾ for a plate the dimensions of which are not small compared to the wavelength.

The reradiating properties of the dipoles when the plate is present, is calculated as if the plate was infinite in extent, using the theory of images. The system of dipoles may then be treated along the same lines as in Scientific Report No. 1, but the induced voltage, the mutual impedances, and the determination of the field reradiated from the dipoles have to be changed because of the image.

The induced voltage is still found as described for an arbitrary reflector in Scientific Report No. 1, but now the electromagnetic field vector is changed in such a way that the distance from the dipoles to the plate is involved, according to ordinary reflection theory.

The new values of self- and mutual impedances are found using the method of images, too.

By using the values of the induced voltages and the values of the self- and mutual impedances thus found, the system of equations (23) of Scientific Report No. 1 will give the new values of the currents on the antenna elements when the presence of the plate is taken into account.

After that the reradiation pattern from the dipole reflector itself mounted in a distance h above the plate is calculated with the above-mentioned currents on the dipoles. The final reradiation pattern of the dipoles is found using the theory of antenna arrays on the array consisting of two parallel Van Atta reflectors in free space with the distance $2h$, where h is the

distance between the dipoles and the plate.

The total field reradiated from the reflector system is found by adding the field reradiated from the dipoles and the field reflected from the conducting plate.

Using the above-mentioned theory a computer program has been developed and the numerical results have been compared with results obtained by Sharp from experimental investigations of a 16 element square Van Atta reflector. The computed back-scattering cross section shows a good agreement with the results measured for the experimental reflector as shown in fig. 5.

Furthermore, a series of computations has been performed in order to examine the changes in the back-scattering cross section due to changing of the parameters of the reflector. The parameters are the number (N) of elements, the length (a) and characteristic impedance (Z_0) of the transmission lines, the distance (d) between adjacent dipole elements, and the distance (h) from the elements to the plate.

The most important results obtained by this numerical investigation is:

- (1) that the back-scattering cross section becomes larger if a distance of 0.35 wavelengths from the dipoles to the plate is used instead of 0.25 wavelengths as used by Sharp,
- (2) that the shape of the curves of back-scattering becomes more irregular when more elements are used in the reflector but the level of back-scattering is increased,
- (3) that the shape of the curves of back-scattering becomes more smooth when a mismatch between the elements and the transmission lines is introduced in such a way that the characteristic impedance of the line is larger than the self-impedance of the dipole. However, the magnitude of the back-scattered energy is then decreasing more rapidly for oblique directions of incidence,
- (4) that for certain lengths of the transmission lines the back-scattering in the direction normal to the reflector tends to zero. This is in accordance with the results obtained for the four element linear reflector mentioned in section 4.

In Scientific Report No. 5 a great number of numerical results obtained by the parameter variation is given as curves of the back-scattering cross section.

6. BANDWIDTH PROPERTIES OF THE SQUARE

VAN ATTA REFLECTOR

The final work on the contract deals with a computation of the bandwidth properties of a 16 element square Van Atta reflector similar to the experimental model used by Sharp.

Part of this work is an investigation of the influence on the back-scattering properties of the reflector of a change in the length of the dipole elements. In the previous investigations on this contract only reflectors consisting of half-wave dipoles have been investigated. The examination of the influence of the length of dipoles is carried out along the same lines as the examination of the influence of the other parameters of the reflector as described in Scientific Report No. 5. This means that the length of the dipoles is changed while keeping all other parameters of the reflector fixed with values corresponding to the dimensions of the experimental reflector used by Sharp.

The results of this investigation is shown in fig. 6 of this report. From this it turns out as expected that the back-scattered energy decreases both when the dipole length is less than and greater than half a wavelength. This is due to the fact that the matched half-wave dipole has optimal reradiating properties. However, when the dipoles are less than half a wavelength the mutual coupling between the dipoles decrease and a better retro-directive effect of the reflector is obtained. This effect is further strengthened because of the mismatch between the dipole and the transmission line, the self resistance R_A of the dipole being less than the characteristic impedance Z_0 of the line. This is in accordance with the results of the investigation of the four element linear reflector where it was found that the deviation from retrodirective effect decreases when the factor R_A/Z_0 decreases.

The bandwidth properties of the 16 element square Van Atta reflector similar to the experimental reflector used by Sharp with or without the conducting plate has been computed and the results are shown in figs. 7a and 8a. The curves show the back-scattered energy for different angles of incidence as a function of λ/λ_0 , where λ_0 is the wavelength corresponding to the center frequency.

For increasing values of λ/λ_0 larger than 1.0 the back-scattered energy from the reflector without a conducting plate decreases in a regular manner for all angles of incidence.

For values of $\lambda/\lambda_0 < 1.0$ the shape of the curves is very irregular and gives no information at all. When the conducting plate is taken into account the curves for $\lambda/\lambda_0 > 1.0$ are almost as regular as when the plate is not present but the level of back-scattered energy is higher for directions of incidence near normal incidence (0°) and lower for directions of incidence near broadside (90°). For $\lambda/\lambda_0 < 1.0$ the curves are just as irregular as in the case where the plate is not present.

However, from the results of this investigation it turned out that the retrodirective effect of this 16 element reflector is essentially improved when $\lambda/\lambda_0 = 1.3$. This corresponds to a 16 element square Van Atta reflector with the following parameter values:

length of dipoles	0.385 λ
radius of dipoles	0.0115 λ
distance between dipoles	0.462 λ
distance from dipoles to plate	0.192 λ
length of transmission lines	0.315 λ
characteristic impedance of transmission lines	73.0 ohms

A measure of retrodirective effect is given in figs. 7b and 8b as

$$M = \frac{\text{number of examined angles of incidence giving retrodirective effect}}{\text{total number of angles of incidence examined}} \cdot 100\%$$

For the reflector without a plate M obtains its maximum 90% for the above-mentioned wavelength $\lambda = 1.3\lambda_0$ of the incident plane wave. When the conducting plate is present the optimum value of M is 50% and this value is obtained in the whole range from $\lambda = 1.1\lambda_0$ to $\lambda = 1.5\lambda_0$. The reason for the decrement of M in the second case is that the retrodirective effect is reduced when the direction of incidence turns towards broadside because of interference with the field scattered from the conducting plate.

However it is obvious in both cases that a Van Atta reflector with a better retrodirective effect than the effect measured by Sharp may be obtained from the same physical reflector at the expense of the reradiated energy if the reflector is used at lower frequencies than the center-frequency.

7. CONCLUSION

In this report a survey of the investigations performed under Contract AF 61(052)-794 "Reflector Array" has been given including the results of the final work on the contract which have not been described in any previous report.

By comparing the results obtained with the objectives of the contract as stated in section 1 of this report it is seen that the first period of the text may be covered by Scientific Reports Nos. 1, 2, 4, and 5. The second period is in part covered by Scientific Report No. 4. The linear and two-dimensional array mentioned in the third period has been investigated while the other configurations mentioned in this period have not been dealt with. The fourth and last period is covered in Scientific Report No. 3 for a linear array.

Further an investigation of the bandwidth properties of a square reflector has been carried out.

Using the theory explained in the reports issued under this contract it should be possible to investigate other types of Van Atta reflectors. This might be as well reflectors consisting of dipoles mounted in two- or three-dimensional arrays for example over conducting cylinders or spheres, as reflectors consisting of other types of antenna elements such as horns, paraboloid antennas, crossed dipoles or monopoles.

Probably results which compare better with the experimental results measured by Sharp may be obtained using another theory for the field reflected from the conducting plate than the physical optics theory used in Scientific Report No. 5. This theory may be Keller's geometrical theory of optics which, in contrast to the theory used, will take into account the scattering of the incident field about the edges of the conducting plate.

8. COMPUTER PROGRAM

In Appendix 1 a copy of the computer program developed for the numerical investigation of square Van Atta reflectors with or without a conducting plate is printed.

The program is written in FORTRAN IV and the computer used in an IBM 7090 run by the Northern Europe University Computing Center, Technical University of Denmark.

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APPENDIX 1

\$IBFTC ATTA DECK

ATTA - EFN SOURCE STATEMENT - IFN(S) -

CALCULATION OF RERADIATION PATTERN OF SQUARE VAN ATTA REFLECTOR
WITH AND WITHOUT A CONDUCTING PLATE

THE INPUT PARAMETERS ARE
COMMON INPUT

I = NUMBER OF ELEMENTS IN EACH ROW
J = 1 IF CONDUCTING PLATE IS TAKEN INTO ACCOUNT, ELSE J = 0
KPO = ANGLE OF INCIDENCE WILL BE CHOSEN AS $KPO + P \cdot 10$ (DEGREES)
MCB = 0 IF RESULTS IN DECIBELS, 1 IF RESULTS IN DIRECT
NUMERICAL VALUES AND 2 IF BOTH CASES ARE WANTED

C = DISTANCE BETWEEN ELEMENTS (IN WAVELENGTHS)
R = DISTANCE FROM ELEMENTS TO PLATE (IN WAVELENGTHS)
B = LENGTH OF TRANSMISSION LINES (IN WAVELENGTHS)
ZO = CHARACTERISTIC IMPEDANCE OF TRANSMISSION LINES (IN OHMS)
RAD = DIPOLE RADII (IN WAVELENGTHS)
DLE = LENGTH OF DIPOLES (IN WAVELENGTHS)

W = ANGLE OF POLARIZATION OF INCIDENT WAVE (IN DEGREES)
FII = ANGLE OF INCIDENCE OF INCOMING PLANE WAVE (IN DEGREES)

LFI = ANGLE OF REFLECTION (IN DEGREES)
MA = 0 IF NEW CALCULATIONS ARE WANTED, ELSE MA = 1
PW = 1 IF BANDWIDTH CALCULATIONS ARE WANTED, ELSE PW
= 0 IF PARAMETER VARIATION IS WANTED

IF BW = 1 (BANDWIDTH CALCULATION) USE FOLLOWING INPUT

MDIP = 0 IF THE PHYSICAL LENGTH OF DIPOLES HAS TO BE UNCHANGED
ELSE MDIP = 1
MD = 0 IF THE PHYSICAL DISTANCE $D \cdot \lambda_{MDAO}$ HAS TO BE UNCHANGED
ELSE MD = 1
MR = 0 IF THE PHYSICAL DISTANCE $R \cdot \lambda_{MDAO}$ HAS TO BE UNCHANGED
ELSE MR = 1
MB = 0 IF THE PHYSICAL LENGTH $B \cdot \lambda_{MDAO}$ HAS TO BE UNCHANGED
ELSE MB = 1

DIFF = HALF THE RANGE OVER WHICH THE FACTOR $\lambda_{MDAO}/\lambda_{MDAO}$
IS VARIED
RATIO = THE STEPS IN WHICH THE FACTOR $\lambda_{MDAO}/\lambda_{MDAO}$ IS VARIED
(λ_{MDAO} = FREQUENCY, λ_{MDAO} = CENTER FREQUENCY)

014CC2

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- EFN SOURCE STATEMENT - IFN(S) -

```

C IF BW = C (PARAMETER VARIATION) USE THE FOLLOWING INPUT
C
C   DI  = STEPS IN CHANGING I DURING PARAMETER VARIATION
C   DD  = STEPS IN CHANGING D DURING PARAMETER VARIATION
C   DR  = STEPS IN CHANGING R DURING PARAMETER VARIATION
C   DDB = STEPS IN CHANGING B DURING PARAMETER VARIATION
C   DZ  = STEPS IN CHANGING ZC DURING PARAMETER VARIATION
C   DDLE = STEPS IN CHANGING DLE DURING PARAMETER VARIATION
C
C   IMX = MAXIMUM VALUE OF I IN PARAMETER VARIATION
C   DMX = MAXIMUM VALUE OF D IN PARAMETER VARIATION
C   RMX = MAXIMUM VALUE OF R IN PARAMETER VARIATION
C   BMX = MAXIMUM VALUE OF B IN PARAMETER VARIATION
C   ZMX = MAXIMUM VALUE OF ZC IN PARAMETER VARIATION
C   CLMX = MAXIMUM VALUE OF DLE IN PARAMETER VARIATION
C
C
C   DIMENSION A(72,72),C(72),X(72),U(72,10)
C   DIMENSION H(36,6),Q(2,10),S(5),V(36,10),IB(6)
C   DIMENSION G(19,10),TETA(10),PHI(10),VV(10),T(19),F(19)
C   INTEGER BW,DI,PAGINA
369  FORMAT(3I5)
301  FORMAT (4I5)
302  FORMAT(4F9.3)
303  FORMAT(2F9.3)
304  FORMAT(3F9.3)
305  FORMAT(6F9.3)
306  FORMAT(15,5F9.3)
100  READ(5,301) I,J,KPQ,MDB
      READ(5,305) D,R,B,ZC,RAC,DLE
      READ(5,303) W,FII
      READ(5,369) LFI,MA,BW
C
C   IF(BW.EQ.1) READ(5,301) MDIP,MD,MR,MB
C   IF(BW.EQ.1) READ(5,303) DIFF,RATIC
C
C   IF(BW.EQ.C) READ(5,306) DI,DD,DR,DCB,DZ,DDLE
C   IF(BW.EQ.C) READ(5,306) IMX,DMX,RMX,BMX,ZMX,CLMX
C
C   IF(BW.EQ.C) DIFF = 0.
C   IF(BW.EQ.C) MDIP = 1
C   K=10
C   LTETA=10
C   FAKTOR=1.C-DIFF
C   PAGINA = C
C   IST=1
C   CST=C
C   RST=R
C   BST=B
C   ZST=ZC
C   CLST=DLE
C
C   IF(BW.EQ.1) GO TO 502
C   IF(BW.EQ.C) GO TO 503
501  I=1 +DI

```

014C02

ATTA

- EFN SOURCE STATEMENT - IFN(S) -

```

GO TO 515
510 C=C +CD
GO TO 515
512 B=B +CCB
GO TO 515
513 ZC=ZC+DZ
GO TO 515
511 R=R +CR
GO TO 515
514 CLE=CLE+CCLE
C
515 IF(BW.EQ.C) GO TO 503
502 IF(MC.EQ.C) C=C/FAKTOR
IF(MR.EQ.C) R=R/FAKTOR
IF(MB.EQ.C) B=B/FAKTOR
503 C=C*6.2831853
R=R*6.2831853
B=B*6.2831853
PAGINA = PAGINA + 1
N=I**2
M=I/2
IF(M*2-I) 11,12,11
11 LE=1
GO TO 2
12 LE=0
2 MM=(I-LE)/2+5/100
CO 13 M=1,MM
13 S(M)=(FLOAT(I)/2.-FLCAT(M)+C.5)*D
LL=(I-LE+1)/2
CO 15 L=1,LL
CO 15 M=1,I
L1=(L-1)*I+M
F(L1,2)= S(L)
L2=I-L+1+(M-1)*I
F(L2,1)= S(L)
L3=N-I+M-(L-1)*I
F(L3,2)=-S(L)
L4=L+(M-1)*I
F(L4,1)=-S(L)
IF(LE) 14,15,14
14 L5=(I-1)/2+(M-1)*I+1
F(L5,1)=C.C
L6=(I*(I-1))/2+M
F(L6,2)=C.C
15 CONTINUE
II=I**2
FAKT1 = 1.CCCCCC
IF(MCIP.EQ.0) FAKT1 = FAKTOR
CALL BETA(II,J,ZC,R,H,FAKT1,RAD,CLE)
TAL=0.
OPQ=FLOAT(KPQ)/10.
TAL=TAL+1.
CLE=CLE/FAKT1
CO 24 M=1,K
YM = M
VV(M)=W

```


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- EFN SOURCE STATEMENT - IFN(S) -

```

PHII(M)=FII
XM = YM-OPQ
XK=K-1
TETAI(M)=(XM*SC.)/XK
IF(ABS(TETAI(M))-90.) 22,21,22
21 IF(ABS(PHII(M))-90.) 22,25,22
22 TETI=TETAI(M)*C.C174532525
Z=0.C174532525
OP=(COS(3.14159265*DLE*SIN(TETI)*SIN(FII*Z))-CCS(3.14159265*CLE))*
1 (SIN(W*Z)*CCS(FII*Z)-COS(W*Z)*SIN(FII*Z)*CCS(TETI))/(SQRT(1.-
2 SIN(TETI)**2*SIN(FII*Z)**2)*SIN(3.14159265*CLE))

```

```

Q(1,M)=P*(1.-FLOAT(J)*CCS(2.*R*CCS(TETI)))
Q(2,M)=-P*FLOAT(J)*SIN(2.*R*CCS(TETI))
DO 23 L=1,N
23 QV(L,M)=-P(L,1)*COS(FII*Z)*SIN(TETI)
1 -H(L,2)*SIN(FII*Z)*SIN(TETI)
24 CONTINUE
GO TO 26
25 K=K-1

```

C
C
C

NOW THE ALXILIARY ARRAY HAS BEEN CALCULATED

C
C
C
C
C

```

26 CALL MAT1(LE,I,N,J,B,IB,A,H)
IF(LE-1) 110,41,110
41 CALL MAT2(LE,I,N,J,B,A,H)
110 CONTINUE
DO 50 L=1,K
IF(LE-1) 49,48,49
48 N1=(N+1)/2
N2=(3*N+1)/2
C(N1) =(Q(1,L)*COS(V(N1,L))-Q(2,L)*SIN(V(N1,L)))*CCS(B)
C(N2) =(Q(1,L)*SIN(V(N1,L))+Q(2,L)*CCS(V(N1,L)))*CCS(B)
49 MM=(N-LE)/2
NL=2*N+L
DO 800 M=1,MM
NM1=N+1-M
MN=N+M
M1=M+(N+LE)/2
M2=M+(3*N+LE)/2
RVM = Q(1,L)*CCS(V(M,L))-Q(2,L)*SIN(V(M,L))
CVM = Q(1,L)*SIN(V(M,L))+Q(2,L)*CCS(V(M,L))
RVM1 = Q(1,L)*COS(V(NM1,L))-Q(2,L)*SIN(V(NM1,L))
CVM1 = Q(1,L)*SIN(V(NM1,L))+Q(2,L)*CCS(V(NM1,L))
C(M) = (RVM-RVM1)*COS(B*C.5)
C(M1) = (RVM+RVM1)*SIN(B*C.5)
C(MN) = (CVM-CVM1)*COS(B*C.5)
800 C(M2) = (CVM+CVM1)*SIN(B*C.5)

```

C
C

NOW THE MATRIX IS FILLED UP AND THE SOLVING OF THE EQUATIONS WILL START

```

IF(L-1) 700,600,700
600 CALL SOLVE(2*N,A,C,1,0.01,5,X,IT)

```

014C02
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- EFN SOURCE STATEMENT - IFN(S) -

GO TO 52
700 CALL SOLVE(2*N,A,C,2,C.01,5,X,IT)
52 N2= 2*N
DO 51 KM=1,N2
51 U(KM,L)=X(KM)
50 CONTINUE

C
C NOW THE RERADIATION PATTERN WILL BE CALCULATED
C

IF(IT) 94,55,55
55 KLM = 0
210 IF(KLM) 230,220,230
220 FI=LFI+180
KT= 90
GO TO 240
230 FI=LFI
KT=100
240 OCALL PATT(KLM,KT,K,I,N,J,D,R,W,ZC,FII,LTETA,FI,TETAI,H,U,G,T,F,CLE
1)
KLM=KLM+1
IF(KLM-1) 281,210,281
281 WRITE (6,300)
WRITE (6,71) N,PAGINA
WRITE (6,72)
WRITE (6,74)
IF(J.EQ.1) WRITE(6,76)
IF(BW.EQ.C) GO TO 310
WRITE(6,200) FAKTOR
IF(MDIP.EQ.0) WRITE(6,333)
IF(MD.EQ.C) WRITE(6,444)
IF(MR.EQ.C) WRITE(6,555)
IF(MB.EQ.0) WRITE(6,666)
310 IF(MDB.NE.C) GO TO 120
DO 111 MK = 1,19
DO 111 KK = 1,K
111 G(MK,KK) = DB(G(MK,KK))
120 AA = B/6.28318531
DPI=D/6.28318531
77 WRITE(6,177) DPI
WRITE(6,777) RAD
WRITE(6,688) DLE
WRITE (6,175) AA
HH=R/6.28318531
IF(J-1) 180,178,180
178 WRITE (6,179) HH
180 WRITE(6,181) Z0
IF(MDB.NE.C) WRITE(6,30)
WRITE (6,78) (TETAI(M),M=1,K)
WRITE (6,79) (PHI1(M),M=1,K)
WRITE (6,80) (VV(M),M=1,K)
WRITE (6,81)
C
DO 84 MK=1,19
84 WRITE (6,90) T(MK),F(MK),(G(MK,M),M=1,K)
IF(MDB.NE.2) GO TO 96
WRITE(6,30)

014CC2

ATTA

- EFN SOURCE STATEMENT - IFA(S) -

```

CO 112 MK = 1,19
CO 112 KK = 1,K
112 G(MK,KK) = DB(G(MK,KK))
CO 113 MK = 1,19
113 WRITE(6,9C) T(MK),F(MK),(G(MK,M),M=1,K)
GO TO 96
94 WRITE(6,95)
96 C=CPI
B=AA
R=FH
IF(BW.EQ.1) GO TO 504
IF(DI.NE.C.AND.I.LT.IMX) GC TC 501
DI=0
I=IST
IF(DC.NE.O..AND.D.LT.DMX) GO TO 510
DD=C.
C=CST
IF(CCB.NE.O..AND.B.LT.BMX) GC TC 512
CDB=0.
B=BST
IF(DZ.NE.C..AND.ZO.LT.ZMX) GC TC 513
CZ=0.
ZO=ZST
IF(DR.NE.C..AND.R.LT.RMX) GC TC 511
CR=0.
R=RST
IF(DDLE.NE.O..AND.DLE.LT.DLMX) GO TO 514
DDLE=0.
DLE=DLST
C
504 IF(BW.EQ.O) GO TO 505
C
IF(MD.EQ.C) D=D*FAKTOR
IF(MR.EQ.C) R=R*FAKTOR
IF(MB.EQ.O) B=B*FAKTOR
DLE=DLE*FAKT1
FAKTOR=FAKTOR+RATIO
IF(FAKTOR.LT.1.C+DIFF) GC TO 502
505 IF(MA) 93,100,93
93 CONTINUE
300 FORMAT (1H1)
71 OFORMAT(31H SQUARE VAN ATTA REFLECTOR WITH,I3,10H ELEMENTS,50X,4P
1AGE,I4)
72 OFORMAT(102H-CALCULATED RERADIATION PATTERN *G* IN THE DIRECTION (I
1VEN BY THE ANGLES TETA AND FI FOR VARIOUS CASES)
74 OFORMAT(67H OF INCIDENCE AND POLARIZATION -GIVEN BY THE ANGLES TETA
1I,FI AND V)
76 OFORMAT(48H WHERE THE CONDUCTING PLATE IS TAKEN INTO ACCOUNT)
200 OFORMAT(45H BANDWIDTH CALCULATIONS FOR LAMBDA/LAMBDA0 =,F8.3,4P AN
1E)
333 OFORMAT(62H NO VARIATION OF PHYSICAL DIPOLE LENGTH (LENGTH= 0.5* LA
1MBDACC))
444 OFORMAT(50H NO VARIATION OF PHYSICAL DISTANCE BETWEEN DIPOLES)
555 OFORMAT(43H NO VARIATION OF PHYSICAL HEIGHT OVER PLATE)
666 OFORMAT(41H NO VARIATION OF PHYSICAL LENGTH OF LINES)
177 OFORMAT (27H DISTANCE BETWEEN ELEMENTS=,F10.2,13H WAVELENGTHS)

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014C02

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- EFN SOURCE STATEMENT - IFN(S) -

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777 FORMAT(27H DIPOLE RADIUS = ,F11.3,12H WAVELENGTHS)
888 FORMAT(27H DIPOLE LENGTH = ,F11.3,12H WAVELENGTHS)
175 FORMAT (29H LENGTH OF TRANSMISSION LINES=,F8.2,13H WAVELENGTHS)
179 FORMAT(33H DISTANCE FROM ELEMENTS TO PLATE=,F4.2,13H WAVELENGTHS)
181 OFORMAT(48H CHARACTERISTIC IMPEDANCE OF TRANSMISSION LINES=,F5.2,6H
1 OHMS)
30 FORMAT(48H RERADIATION PATTERN VALUES MEASURED IN DECIBELS)
78 FORMAT(9H-TETA I= ,2X,10F9.1)
79 FORMAT(9H FII = ,2X,10F9.1)
80 FORMAT(9H V = ,2X,10F9.1)
81 FORMAT(11H TETA FI)
90 FORMAT(1H ,F4.1,F6.0,10F9.2)
95 FORMAT(42H1***THE SYSTEM OF EQUATIONS IS SINGULAR***)
STOP
END

```

\$IBFTC DECI DECK

REAL FUNCTION DB(X)

```

C
C * DB * CONVERTS THE RERADIATED ENERGY INTO DECIBELS
C
DB = -900.0
XI = 10000000.0 * X
KX = IFIX(XI)
IF(KX.NE.0) GO TO 1
RETURN
1 DB = 10.0 * ALOG10(X)
RETURN
END

```

014002

\$IBFTC SUB1 CECK

SUB1 - EFN SOURCE STATEMENT - IFN(S) -

```

SUBROUTINE MAT1(LE,I,N,J,B,IB,A,H)
C
C *MAT1* COMPUTES THE FIRST PART OF THE MATRIX EQUATION
C
  DIMENSION IB(6),A(72,72),H(36,6)
  MM=(N-LE)/2
  DO 32 M=1,MM
  DO 30 L=1,I
  MX=M-L+I
  IF(100-MX-5) 27,30,30
  27 IF(100-(MX-LE)+50*I-5) 28,29,29
  28 LL=N-2*M+2
  GO TO 31
  29 LL=2*MX+N-2*L+I+2*I
  GO TO 31
  30 CONTINUE
C
  31 Y=((H(1,3)-FLOAT(J)*H(1,5)-H(LL,3)+FLOAT(J)*H(LL,5))*CCS(B*0.5)
  A(M,2*M-1)= Y
  MN=M+N
  A(MN,2*M)= Y
  NM1=2*(N-M)+1
  NM2=2*(N-M)+2
  ML1=M+(N+LE)/2
  ML2=M+(3*N+LE)/2
  A(M,NM1)= -Y
  A(MN,NM2)= -Y
C
  OY=-(H(1,4)-FLCAT(J)*H(1,6)-H(LL,4)+FLCAT(J)*H(LL,6))*CCS(B*0.5)
  1 +SIN(B*0.5)
  A(M,2*M)= Y
  A(MN,NM1)= Y
  A(M,NM2)= -Y
  A(MN,2*M-1)= -Y
C
  Y=(H(1,3)-FLOAT(J)*H(1,5)+H(LL,3)-FLCAT(J)*H(LL,5))*SIN(B*0.5)
  A(ML1,2*M-1)= Y
  A(ML1,NM1)= Y
  A(ML2,2*M)= Y
  A(ML2,NM2)= Y
C
  OY=-(H(1,4)-FLOAT(J)*H(1,6)+H(LL,4)-FLCAT(J)*H(LL,6))*SIN(B*0.5)
  1 -COS(B*0.5)
  A(ML1,2*M)= Y
  A(ML1,NM2)= Y
  A(ML2,2*M-1)= -Y
  A(ML2,NM1)= -Y

```

014002

SUB1

- EFN SCLRC STATEMENT - IFN(S) -

32 CONTINUE

MM=(N-3*LE)/2

LL=(N-LE)/2

CO 40 M=1,MM

ML=M+1

IF(LL-ML) 559,668,888

888 CO 40 L=ML,LL

K1=C

L2= L-M

CO 35 M1=M,L,L2

K1=K1+1

CO 34 L1=1,LL

MX=M1-L1*1

IF(2*MX-1) 33,34,34

33 IB(K1)=MX+1-1

K1=K1+1

IB(K1)=L1-1

GO TO 35

34 CONTINUE

35 CONTINUE

IB(5)=1-1-IB(3)

IB(6)=1-1-IB(4)

IF(1B(1)-1B(3)) 37,37,36

36 LL1= 1*(1B(4)-1B(2))-1B(3)+1B(1)+1

GO TO 113

37 LL1= 1*(1B(4)-1B(2))+1B(3)-1B(1)+1

113 IF(1B(1)-1B(5)) 39,39,38

38 LL2=1*(1B(6)-1B(2))-1B(5)+1B(1)+1

GO TO 114

39 LL2= 1*(1B(6)-1B(2))+1B(5)-1B(1)+1

114 OY=(H(LL1,3)-FLOAT(J)*H(LL1,5)-H(LL2,3)+FLCAT(J)*H(LL2,5))

1 *COS(B*C.5)

A(L,2*M-1)=Y

A(M,2*L-1)=Y

NL=N+L

NM=N+M

NM1= 2*(N-M)+1

NL1=2*(N-L) +1

NM2=NM1+1

NL2=NL1+1

LE1=L+(N+LE)/2

ME1=M+(N+LE)/2

LE2=L+(3*N+LE)/2

ME2=M+(3*N+LE)/2

A(NM,2*L)= Y

A(NL,2*M)= Y

A(L,NM1)= -Y

A(M,NL1)= -Y

A(NL,NM2)= -Y

A(NM,NL2)= -Y

OY=-(H(LL1,4)-FLOAT(J)*H(LL1,6)-H(LL2,4)+FLCAT(J)*H(LL2,6))

1 *COS(B*C.5)

A(L,2*M)=Y

014C02

SUB1

- EFN SCLRC STATEMENT - IFN(S) -

A(M,2*L)=Y
A(NL,NM1)=Y
A(NM,NL1)=Y
A(L,NM2)=-Y
A(M,NL2)=-Y
A(NL,2*M-1)=-Y
A(NM,2*L-1)=-Y

C

OY=(H(LL1,3)-FLCAT(J)*H(LL1,5)+H(LL2,3)-FLCAT(J)*H(LL2,5))

1 *SIN(B*C.5)
A(LE1,2*M-1)=Y
A(ME1,2*L-1)=Y
A(LE1,NM1)=Y
A(ME1,NL1)=Y
A(LE2,2*M)=Y
A(ME2,2*L)=Y
A(LE2,NM2)=Y
A(ME2,NL2)=Y

C

OY=-(H(LL1,4)-FLOAT(J)*H(LL1,6)+H(LL2,4)-FLCAT(J)*H(LL2,6))

1 *SIN(B*C.5)
A(LE1,2*M)=Y
A(ME1,2*L)=Y
A(LE1,NM2)=Y
A(ME1,NL2)=Y
A(LE2,2*M-1)=-Y
A(ME2,2*L-1)=-Y
A(LE2,NM1)=-Y
A(ME2,NL1)=-Y

40 CONTINUE
999 CONTINUE
RETURN
END

014CC2

\$IBFTC SUB2 DECK

SUB2 - EFN SOURCE STATEMENT - IFN(S) -

```

SUBROUTINE MAT2(LE,I,N,J,B,A,H)
C
C *MAT2* COMPUTES THE SECCND PART CF THE MATRIX EQUATION
C
DIMENSION A(72,72),H(36,6)
Y=(I(1,3)-FLOAT(J)*H(1,5))*CCS(B)
C
NE=(N+LE)/2
NE1=(3*N+LE)/2
A(NE,N)= Y
A(NE1,N+1)= Y
A(NE,N+1)= -(H(1,4)-FLOAT(J)*H(1,6))*CCS(B)+SIN(B)
A(NE1,N)= -A(NE,N+1)
C
MM=(N-LE)/2
DO 47 M=1,MM
LL=(N-LE)/2
DO 43 L=1,LL
MX=M-L+1
IF(2*MX-1) 42,43,43
42 NX=MX+I-1
NY=I-1
GO TO 44
43 CONTINUE
44 IF(2*NX-1) 46,46,45
45 LL3= I*((I-1)/2-NY)+(3-I)/2+NX
GO TO 115
46 LL3= I*((I-1)/2-NY)+(1+I)/2-NX
C
115 Y=(H(LL3,3)-FLOAT(J)*H(LL3,5))*COS(B)
NM=N+M
N1=(N+1)/2
N2=(3*N+1)/2
NM1=2*(N-M)+1
NM2= NM1+1
M2= (3*N+1)/2+M
M1= (N+1)/2+M
A(N1,2*M-1)= Y
A(N1,NM1)= Y
A(N2,NM2)= Y
A(N2,2*M)= Y
C
Y=-(H(LL3,4)-FLOAT(J)*H(LL3,6))*CCS(B)
A(N1,2*M)= Y
A(N1,NM2)= Y
A(N2,2*M-1)= -Y
A(N2,NM2)= -Y

```


014CC2

SUB2

- EFN SOURCE STATEMENT - IFN(S) -

C

```
Y=2.*(H(LL3,3)-FLOAT(J)*H(LL3,5))*SIN(B*0.5)
A(M2,N+1)= Y
A(M1,N)= Y
A(M1,N+1)= -2.*(H(LL3,4)-FLCAT(J)*H(LL3,6))*SIN(B*0.5)
A(M2,N)= -A(M1,N+1)
A(M,N)=0.C
A(M,N+1)=C.C
A(NM,N)= C.C
47 A(NM,N+1)= 0.C
RETURN
END
```

014C02

\$IBFTC SUB3 CECK

SUB3 - EFN SOURCE STATEMENT - IFN(S) -

```

OSUBROUTINE PATT(KLM,KT,K,I,N,J,D,R,W,ZC,FII,LTETA,FI,TETA,H,U,
1          G,T,F,DLE)
C
C      *PATT* COMPUTES THE RERADIATION PATTERN OF THE REFLECTOR
C
      DIMENSION TETA1(10),H(36,6),L(72,10),G(19,10),T(19),F(19)
      PO(X)= SIN(0.5*X*XI*D)/(0.5*X*XI*C)
      MKM=0
      DO 91 NTETA=10,KT,LTETA
      MKM=MKM+1
      MK= KLM* 9+MKM
      NX = (MKM-1)*(KLM-1)+ MKM * KLM
      IF(KLM) 260,250,260
250  MTETA=100-NTETA
      GO TO 270
260  MTETA=NTETA - 10
270  DO 89 M=1,K
      GR=0.
      GI=0.
      Y=3.14159265/180.
      TETA=MTETA
      TETA = TETA1(NX)
      T(MK)=TETA
      F(MK)=FI
      OXK1=120.*3.14159265*(COS(3.14159265*DLE*SIN(TETA*Y)*SIN(FI*Y))
1      -COS(3.14159265*DLE))/(ZC*SQRT(1.-(SIN(TETA*Y)*SIN(FI*Y))**2)*
2      SIN(3.14159265*DLE))

      N2=2*N
      DO 84 L=1,N2,2
      L1=(L+1)/2
      C1=H(L1,1)*SIN(TETA*Y)*COS(FI*Y)+H(L1,2)*SIN(TETA*Y)*SIN(FI*Y)
      AA= U(L,M)
      NL=L+1
      BB= U(NL,M)
      GR=GR+AA*COS(C1)+BB*SIN(C1)
      GI=GI-AA*SIN(C1)+BB*COS(C1)
84  CONTINUE
      IF(J-1) 87,85,87
85  DO 86 L=1,N2,2
      L1=(L+1)/2
      OC1=H(L1,1)*SIN(TETA*Y)*COS(FI*Y)+H(L1,2)*SIN(TETA*Y)*SIN(FI*Y)
1      -2.*R*COS(TETA*Y)
      AA= U(L,M)
      NL=L+1
      BB= U(NL,M)
      AN=N

```

C1400?

SUB3

- EFN SOURCE STATEMENT - IFN(S) -

```

GR=GR-AA*COS(C1)-BB*SIN(C1)
GI=GI+AA*SIN(C1)-BB*CCS(C1)
86 CONTINUE
XI=I
AR=COS(R*(COS(TETAI(M)*Y)+CCS(TETA*Y)))
AI=SIN(R*(COS(TETAI(M)*Y)+CCS(TETA*Y)))
AL=SIN(TETAI(M)*Y)*CCS(FII*Y)+SIN(TETA*Y)*CCS(FI*Y)
BF=SIN(TETAI(M)*Y)*SIN(FII*Y)+SIN(TETA*Y)*SIN(FI*Y)
AL1=AL*100000.
BE1=BE*100000.
IF(IFIX(AL1)) 290,291,290
291 SAL=1.
GO TO 292
290 SAL = PO(AL)
292 IF(IFIX(BE1)) 294,293,294
293 SBE=1.
GO TO 295
294 SBE = PO(BE)
295 XK2= C.5*SAL*SBE*AN*D**2
TP=COS(W*Y)*COS(FII*Y)+SIN(W*Y)*CCS(TETAI(M)*Y)*SIN(FII*Y)
TT=-COS(W*Y)*SIN(FII*Y)+SIN(W*Y)*CCS(TETAI(M)*Y)*CCS(FII*Y)
FT=-TP*CCS(FI*Y)*CCS(TETA*Y)+TT*SIN(FI*Y)*CCS(TETA*Y)
FP=TP*SIN(FI*Y)+TT*CCS(FI*Y)
OG(MK,M)= ((XK1*GR*SIN(FI*Y)*CCS(TETA*Y) + XK2*AR*FT)**2
1          +(XK1*GI*SIN(FI*Y)*CCS(TETA*Y) + XK2*AI*FT)**2
2          +(XK1*GR*CCS(FI*Y) + XK2*AR*FP)**2
3          +(XK1*GI*CCS(FI*Y) + XK2*AI*FP)**2)/(3.14159265**3)

87 IF(J-1) EE,EE,EE
88 OG(MK,M)=XK1**2*(GR**2+GI**2)*((SIN(FI*Y))**2*(CCS(TETA*Y))**2
1          +(CCS(FI*Y))**2)/(3.14159265**3)
89 CONTINUE
91 CONTINUE
RETURN
END

```

014002

\$IBFTC IMPZ DECK

IMPZ - EFN SOURCE STATEMENT - IFN(S) -

```

SUBROUTINE BETA(II,J,ZC,R,H,FAKTOR,RAD,DLE)
C
C *BETA* BUILDS UP AN ARRAY OF SELF- AND MUTUAL IMPEDANCES
C
  DIMENSION F(36,6)
  CALL SELF(DLE/FAKTOR,RAD,RS,XS)
  DO 20 M=1,II
    IF(M-1) 17,16,17
  16 F(1,3) = RS/ZC
    F(1,4) = -XS/ZC
    GO TO 18
  17 OZ= AMPE(DLE/FAKTOR,DLE/FAKTOR,
    1      ABS(C.1591549*(H(1,1)-H(M,1))),
    1      ABS(C.1591549*(H(1,2)-H(M,2))),0.0,0.0,1,0.001,10)
    F(M,3)=Z/ZO
    OZ= AMPE(DLE/FAKTOR,DLE/FAKTOR,
    1      ABS(C.1591549*(H(1,1)-H(M,1))),
    1      ABS(C.1591549*(H(1,2)-H(M,2))),0.0,0.0,0,0.001,10)
    F(M,4)=Z/ZO
  18 IF(J) 19,20,19
  19 OZ= AMPE(DLE/FAKTOR,DLE/FAKTOR,
    1      0.1591549*SQRT((H(1,1)-H(M,1))*2+4.*(R**2)),
    1      ABS(C.1591549*(H(1,2)-H(M,2))),0.0,0.0,1,0.001,10)
    F(M,5)=Z/ZC
    OZ= AMPE(DLE/FAKTOR,DLE/FAKTOR,
    1      C.1591549*SQRT((H(1,1)-H(M,1))*2+4.*(R**2)),
    1      ABS(C.1591549*(H(1,2)-H(M,2))),0.0,0.0,0,0.001,10)
    F(M,6)=Z/ZO
  20 CONTINUE
    RETURN
    END

```

014002

\$IBFTC XCAL DECK

XCAL - EFN SOURCE STATEMENT - IFN(S) -

```

SUBROUTINE SELF(H,A,R,X)
C
C *SELF* COMPUTES THE SELF IMPEDANCE OF THE CIRCLES
C
COMPLEX CSINT,CINSI
PI=3.14159265
AA=SQRT(2.)*2.*PI*A
FH=PI*F
CSINT=CINSI(2.*FH)
CI2= 0.577215665+ALOG(2.*HH)-REAL(CSINT)
S2= REAL(CSINT)
SI2= AIMAG(CSINT)
CSINT= CINSI(4.*FH)
CI4= 0.577215665+ALOG(4.*HH)-REAL(CSINT)
S4= REAL(CSINT)
SI4= AIMAG(CSINT)
X= -29.997925*(SIN(2.*HH)*1-C.577215665+ALOG(HH/(AA**2))+2.*CI2
C      -CI4) - COS(2.*HH)*(2.*SI2-SI4)-2.*SI2)
R= 29.997925*((2.+2.*COS(2.*HH))*S2-COS(2.*HH)*S4
C      -2.*SIN(2.*HH)*SI2+SIN(2.*HH)*SI4)
X=X/(SIN(FH)**2)
R=R/(SIN(FH)**2)
12 CONTINUE
13 CONTINUE
RETURN
END

```

014C02

\$IBFTC CINSI DECK

CINSI - EFN SOURCE STATEMENT - IFN(S) -

```
      COMPLEX FUNCTION CINSI(X)
C CINSI COMPUTES AS ITS REAL PART THE MODIFIED CCSINE INTEGRAL AND AS
C ITS IMAGINARY PART THE SINE INTEGRAL.
      IF (X.GT.1.1) GO TO 20
      C=1.0
      S=1.0
      TC=1.0
      TS=1.0
      Y=X*X
      I=0
10    I=I+1
      TC=-Y*TC/FLOAT(2*I*(2*I-1))
      TS=-Y*TS/FLOAT(2*I*(2*I+1))
      TERMC=TC/FLOAT(2*I)
      TERMS=TS/FLOAT(2*I+1)
      C=C+TERMC
      S=S+TERMS
      EC=ABS(TERMC/C)
      ERROR=1.0E-8
      IF (EC.GT.ERROR) GO TO 10
      CIN=1.0-C
      SI=X*S
      CINSI=CMPLX(CIN,SI)
      RETURN
20    Y=X*X
      F=(Y*(Y*(Y*(Y*(Y+38.027264)+265.187033)+335.677320)+38.102495)/
C (X*(Y*(Y*(Y*(Y*(Y+40.021433)+322.624911)+570.236280)+157.105423))
      G=(Y*(Y*(Y*(Y*(Y+42.242855)+302.757865)+352.018498)+21.821899)/
C (Y*(Y*(Y*(Y*(Y*(Y+48.196927)+482.485984)+1114.978885)+449.690326))
      S=SIN(X)
      C=COS(X)
      CIN=0.577215665+ALOG(X)-F*S+G*C
      SI=1.57079633-F*C-G*S
      CINSI=CMPLX(CIN,SI)
      RETURN
      END
```

C14002

118FTC ZCAL DECK

ZCAL - EFN SOURCE STATEMENT - IFN(S) -

```

REAL FUNCTION AMPE(H1,H2,YC,ZC,TETA,FI,NC,DELTA,MCRD)
C
C *AMPE* COMPUTES THE MUTUAL IMPEDANCES BETWEEN THE CIRCLES BY AN
C INTEGRATION USING ROMBERG'S METHOD
C
C THE REAL PART OF THE MUTUAL IMPEDANCE FOR NC=1 AND
C THE IMAGINARY PART FOR NC=C
C
  DIMENSION TRAP(11)
  STEP= 1/2
  X= -H2/2.
  C1=RX(H1,H2,YC,ZC,TETA,FI,NC,X)
  X=H2/2.
  FC=RX(H1,H2,YC,ZC,TETA,FI,NC,X)
  TRAP(1)=(C1+FC)*STEP/2.
  C1=C.
  DO 203 K=1,MCRD
    SUM=C.
    ERROR=C.
    STEP=STEP/2.
    M=2**K
    MM=M-1
    DO 204 L=1,MM,2
      LL=M-L
      XL=FLOAT(LL)
      XM=FLOAT(M)
      X=XL/XM
      X=X*(-H2/2.)+(1.-X)*H2/2.
      F0=RX(H1,H2,YC,ZC,TETA,FI,NC,X)
      C2=SUM+FC
      IF(ABS(FC)-ABS(SUM)) 206,206,205
205 ERROR=ERROR+SUM-(C2-FC)
      GO TO 207
206 ERROR=ERROR+FC-(C2-SUM)
207 SUM=C2
204 CONTINUE
      TRAP(K+1)=TRAP(K)/2.+(C2+ERROR)*STEP
      P=1.
      KK=10-K
      DO 208 LL=KK,9,1
        L=-(LL-10)
        P=P*4.
208 TRAP(L)=(TRAP(L+1)*P-TRAP(L))/(P-1.)
        C2= TRAP(1)
        IF(K-1) 209,209,211
211 IF(ABS(C2-C1)-DELTA*ABS(C2)) 210,210,209
209 C1=C2

```

C14002

ZCAL

- EFN SOURCE STATEMENT - IFN(S) -

```
203 CONTINUE
    MORC=R-1
210 Q=29.97925/SIN(3.14159265*H1)/SIN(3.14159265*H2)
    IF(NO-1) 201,200,201
200 AMPE= -Q*C2
    GOTO 202
201 AMPE=  Q*C2
202 CONTINUE
    RETURN
    END
```


014002

\$IBFTC INTG DECK

INTG - EFN SOURCE STATEMENT - IFN(S) -

```

REAL FUNCTION RX(H1,H2,YC,ZC,TETA,FI,NC,X)
C
C * RX * COMPUTES THE FUNCTION TO BE INTEGRATED IN ROUTINE AMPE
C
  B1=H1/2.
  B2=H2/2.
  A=COS(TETA)
  B=SIN(TETA)
  C=B*SIN(FI)
  D=B*COS(FI)
  E=2.*COS(3.14159265*H1)
  F=Z0+B1
  G=Z0-B1
  SX=X*C
  SY=X*C
  SZ=X*A
  RO2=SX**2 +(YC+SY)**2
  F=Z0+SZ
  R=SQRT(RO2+H**2 )
  FH=F+SZ
  P=G+SZ
  R1=SQRT(RO2+FH**2 )
  R2=SQRT(RO2+P**2 )
  IF(N0-1) 214,211,214
211 AK=SIN(6.28318531*R1)/R1
  AL=SIN(6.28318531*R2)/R2
  AM=SIN(6.28318531*R)/R
  GO TO 212
214 AK=COS(6.28318531*R1)/R1
  AL=COS(6.28318531*R2)/R2
  AM=COS(6.28318531*R)/R
212 CONTINUE
  IF(Y0) 216,215,218
215 IF(SX) 218,216,218
216 IF(SY) 218,217,218
217 RX=(E*AM-AK-AL)*A*SIN(6.28318531*(B2-ABS(X)))
  GO TO 219
218 ORX=((AK*FH+AL*P-E*AM*H)*(X*(C**2)+Y0*C+X*(C**2))/RC2
  1 +(E*AM-AK-AL)*A)*SIN(6.28318531*(B2-ABS(X)))
219 CONTINUE
  RETURN
  END

```

014C02

\$IBFTC SOLV DECK

SOLV - EFN SOURCE STATEMENT - IFN(S) -

```

SUBROUTINE SOLVE(NN,A,B,IN,EPS,ITMAX,X,IT)
C
CSOLVE  LINEAR EQUATION SOLVER WITH ITERATIVE IMPROVEMENT  VERSION IV
C SOLVES AX=B WHERE A IS NXN MATRIX AND B IS NX1 VECTOR
C IN=
C     1 FOR FIRST ENTRY
C     2 FOR SUBSEQUENT ENTRIES WITH NEW B
C     3 TO RESTORE A AND B
C EPS AND ITMAX ARE PARAMETERS IN THE ITERATION
C IT=
C     -1 IF A IS SINGULAR
C     0 IF NOT CONVERGENT
C     NUMBER OF ITERATIONS IF CONVERGENT
C CALLS MAP SUBROUTINES ILOG2,DCT,SDCT AND CAD
C
C TO MODIFY DIMENSIONS, CHANGE THE NEXT 3 (ACT 2 BUT 3) CARDS.
C
OCIMENSION A(72,72),B(72),X(72),DX(72),R(72),Z(72),RM(72),IRP(72),
IAA(72,72)
MA=72
C MA MUST = DECLARED DIMENSION OF SYSTEM
EQUIVALENCE(R,DX)
GO TO (1000,2000,3000),IN
1000 N=NN
NM1=N-1
NP1=N+1
C
C EQUILIBRATION
C
DO 510 I=1,N
KTOP=ILOG2(A(I,1))
DO 503 J=2,N
503 KTOP=MAXC(KTOP,ILOG2(A(I,J)))
RM(I)=2.C*(-KTOP)
DO 509 J=1,N
509 A(I,J)=A(I,J)*RM(I)
510 CONTINUE
C
C SAVE EQUILIBRATED DATA
C
DO 548 I=1,N
DO 548 J=1,N
548 AA(I,J)=A(I,J)
C
C GAUSSIAN ELIMINATION WITH PARTIAL PIVOTING
C
DO 99 M=1,NM1

```

```

014002
SOLV      - EFN  SOURCE STATEMENT - IFN(S) -

      TOP=ABS (A(N,M))
      IMAX=M
      DO 12 I=M,N
        IF(TOP-ABS (A(I,M)))10,12,12
10        TOP=ABS (A(I,M))
        IMAX=I
12      CONTINUE
        IF(TOP)14,13,14
13      IT=-1
C      *SINGULAR*
      RETURN
14      IRP(M)=IMAX
23      IF(IMAX-M)29,29,24
24      DO 25 J=1,N
        TEMP=A(M,J)
        A(M,J)=A(IMAX,J)
        A(IMAX,J)=TEMP
25      MPI=M+1
29      DO 33 I=MPI,N
        EM=A(I,M)/A(M,M)
        A(I,M)=EM
        IF(EM)31,33,31
31      DO 32 J=MPI,N
32      A(I,J)=A(I,J)-A(M,J)*EM
33      CONTINUE
99      CONTINUE
      IRP(N)=N
      IF (A(N,N))120,113,120
113     IT=-1
      RETURN
120     CONTINUE
C      STORAGE FOR A NOW CONTAINS TRIANGULAR L AND U SO THAT (L+I)*U=A
C
C      DUPLICATE INTERCHANGES IN DATA
C
      DO 229 I=1,N
        IP=IRP(I)
        IF(I-IP)221,229,221
221      DO 222 J=1,N
        TEMP=AA(I,J)
        AA(I,J)=AA(IP,J)
        AA(IP,J)=TEMP
222      CONTINUE
229     CONTINUE
C
C      PROCESS RIGHT HAND SIDE
C
2000    CONTINUE
      DO 601 I=1,N
      601      B(I)=B(I)*RM(I)
      DO 609 I=1,NM1
        IP=IRP(I)
        TEMP=B(I)
        B(I)=B(IP)
        B(IP)=TEMP
      609    CONTINUE
C

```

014C02
SOLV - EFN SOURCE STATEMENT - IFN(S) -

```

C SOLVE FOR FIRST APPROXIMATION TC X
C
199 DO 200 I=1,N
200 Z(I)=-SDOT(I-1,A(I,1),MA,Z(1),1,-B(I))
    DO 201 K=1,N
    I=NP1-K
201 X(I)=-SDOT(N-I,A(I,I+1),MA,X(I+1),1,-Z(I))/A(I,I)
C
C ITERATIVE IMPROVEMENT
C
    IF(ITMAX)370,370,300
300 TOP=0.0
    DO 302 I=1,N
302 TOP=AMAX1(TOP,ABS(X(I)))
    EPSX=EPS*TOP
    DO 369 IT 1,ITMAX
C FIND RESIDUALS
    DO 319 I=1,N
319 R(I)=-DOT(N,AA(I,1),MA,X(1),1,-B(I))
C FIND INCREMENT
    DO 329 I=1,N
329 Z(I)=-SDOT(I-1,A(I,1),MA,Z(1),1,-R(I))
    DO 339 K=1,N
    I=NP1-K
339 DX(I)=-SDOT(N-I,A(I,I+1),MA,DX(I+1),1,-Z(I))/A(I,I)
C INCREMENT AND TEST CONVERGENCE
    TOP=C.0
    DO 342 I=1,N
        TEMP=X(I)
        X(I)=CAD(X(I),DX(I))
        DELX=ABS (X(I)-TEMP)
        TOP=AMAX1(TOP,DELX)
342 CONTINUE
    IF(TOP-EPSX)381,381,369
369 CONTINUE
370 IT=0
381 RETURN
C
C RESTORE A AND B
C
3000 CONTINUE
    DO 709 K=1,N
    I=NP1-K
    IP=IRP(I)
    IF(I-IP)701,709,701
701 TEMP=B(I)
    B(I)=B(IP)
    B(IP)=TEMP
    DO 702 J=1,N
        TEMP=AA(I,J)
        AA(I,J)=AA(IP,J)
        AA(IP,J)=TEMP
702 CONTINUE
709 DO 729 I=1,N
    B(I)=B(I)/RM(I)
    DO 729 J=1,N

```

014C02
SOLV

- EFN SCLRCE STATEMENT - IFN(S) -

A(I,J)=AA(I,J)/RM(I)
729 CONTINUE
RETURN
END

*I8FTC LOG CECK

C INTEGER FUNCTION ILOG2(Z)
C
C *ILOG2* ROUTINE FOR LSE WITH ROUTINE SOLVE
C
ILOG2=C
IF (Z.NE.C.) GO TO 1
RETURN
1 ILOG2=AINTE(3.322C*ALOG10(ABS(Z)))
RETURN
END

C14CC2

\$IDMAP DOT. 84

* DOT AND FRIENDS ROUTINES FOR USE WITH SOLVE
 ENTRY DOT (N,A(1),MA,B(1),MB,C) DOUBLE INNER PRODUCT
 ENTRY SDOT (N,A(1),MA,B(1),MB,C) INNER PRODUCT
 ENTRY ILOG2 (A) FLOATING POINT EXPONENT
 ENTRY DAD (A,B) ADD WITH ROUND

* SNAD MACRO M STORE NEGATIVE OF ADDRESS IN DECREMENT
 SUB =0100000 COMPLEMENT IF POSITIVE
 ALS 18
 STD M
 ENDM SNAD

* DOT SAVE 1,2,4

STZ S
 STZ S+1
 CLA* 8,4 C
 LDQ C+1
 STO C
 CLA* 3,4 N
 TZE NONE SKIP LOOP IF N = 0
 STO N

CLA 4,4 BASE ADDRESS OF A
 PAC ,1 X1=-(BASE OF A)
 CLA* 5,4 MA
 SNAD MA
 CLA 6,4 BASE ADDRESS OF B
 PAC ,2 X2=-(BASE OF B)
 CLA* 7,4 MB
 SNAD MB
 LXA N,4 X4=N
 LOOP LDQ 0,1 A(1)
 FMP 0,2 B(1)
 DFAD S
 DST S

MA TXI **1,1,** (X1)=(X1)+MA
 MB TXI **1,2,** (X2)=(X2)+MB
 TIX LOOP,4,1 END OF MAIN LOOP
 NONE DFAD C

07/13/67

	FRN	
	RETURN	DOT
* SDOT	SAVE	1,2,4
	STZ	S
	CLA*	8,4
	STO	C
	CLA*	3,4
	TZE	SNONE
	STO	N
	CLA	4,4
	PAC	,1
	CLA*	5,4
	SNAD	SMA
	CLA	6,4
	PAC	,2
	CLA*	7,4
	SNAD	SMB
	LXA	N,4
SLOOP	LDQ	0,1
	FMP	0,2
	FAD	S
	STO	S
SMA	TXI	**1,1,**
SMB	TXI	**1,2,**
	TIK	SLOOP,4,1
SNONE	FAD	C
	RETURN	SDOT
* ILOG2	CAL*	3,4
	ANA	=03770000000000
	SUB	=02000000000000
	ARS	27
	TRA	1,4
* DAD	CLA*	3,4
	FAD*	4,4
	FRN	
	TRA	1,4
*		
	EVEN	

- 46 -

07/13/67

C	PZE
	PZE
S	PZE
	PZE
N	PZE
	*LDIR
	*LORG
	END

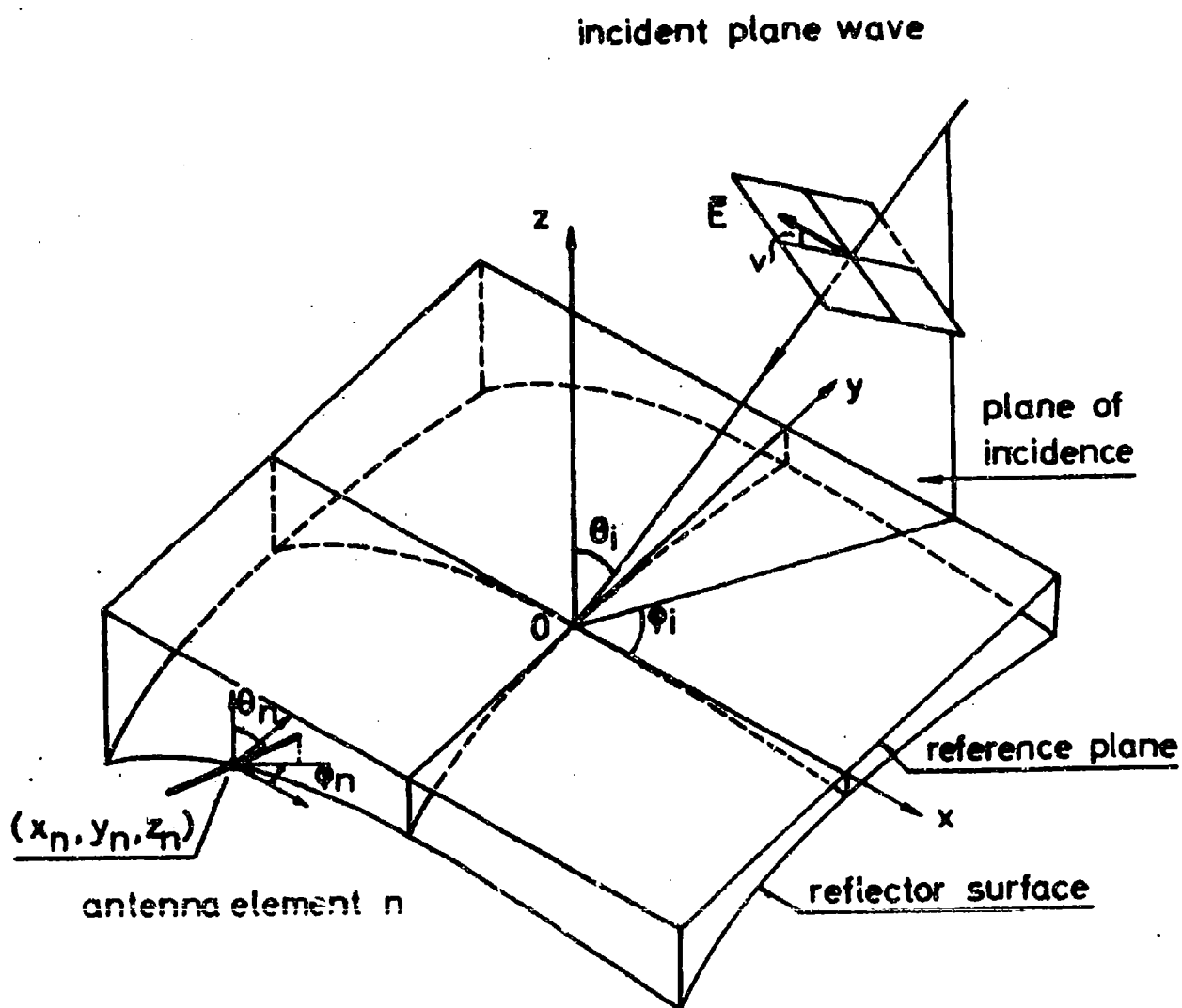


Fig1. Coordinate system for arbitrary Van Atta array.

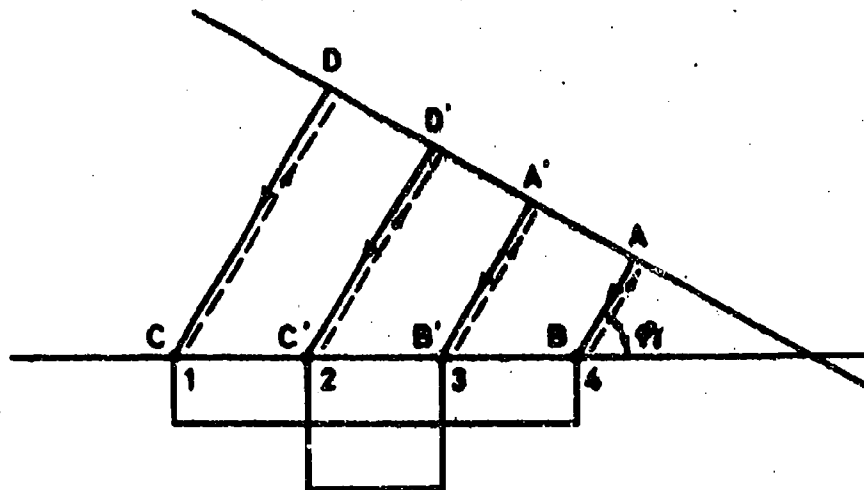


Fig2. The Van Atta principle (retrodirective effect).
The paths ABCD and A'B'C'D' are equal.

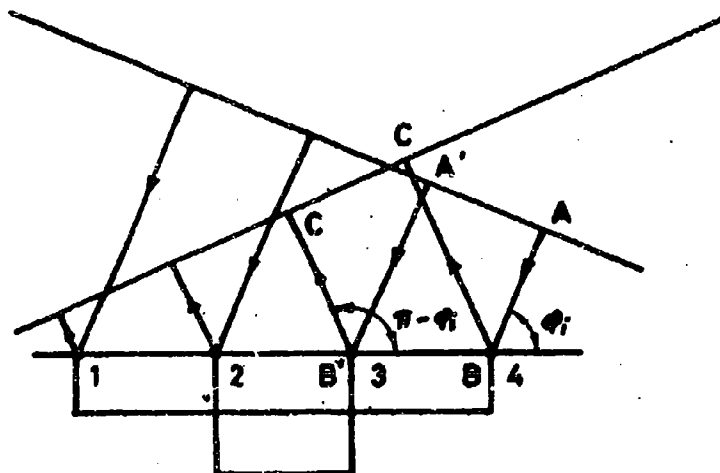


Fig3. The specular reflection (mirror effect).
The paths ABC and A'B'C' are equal.

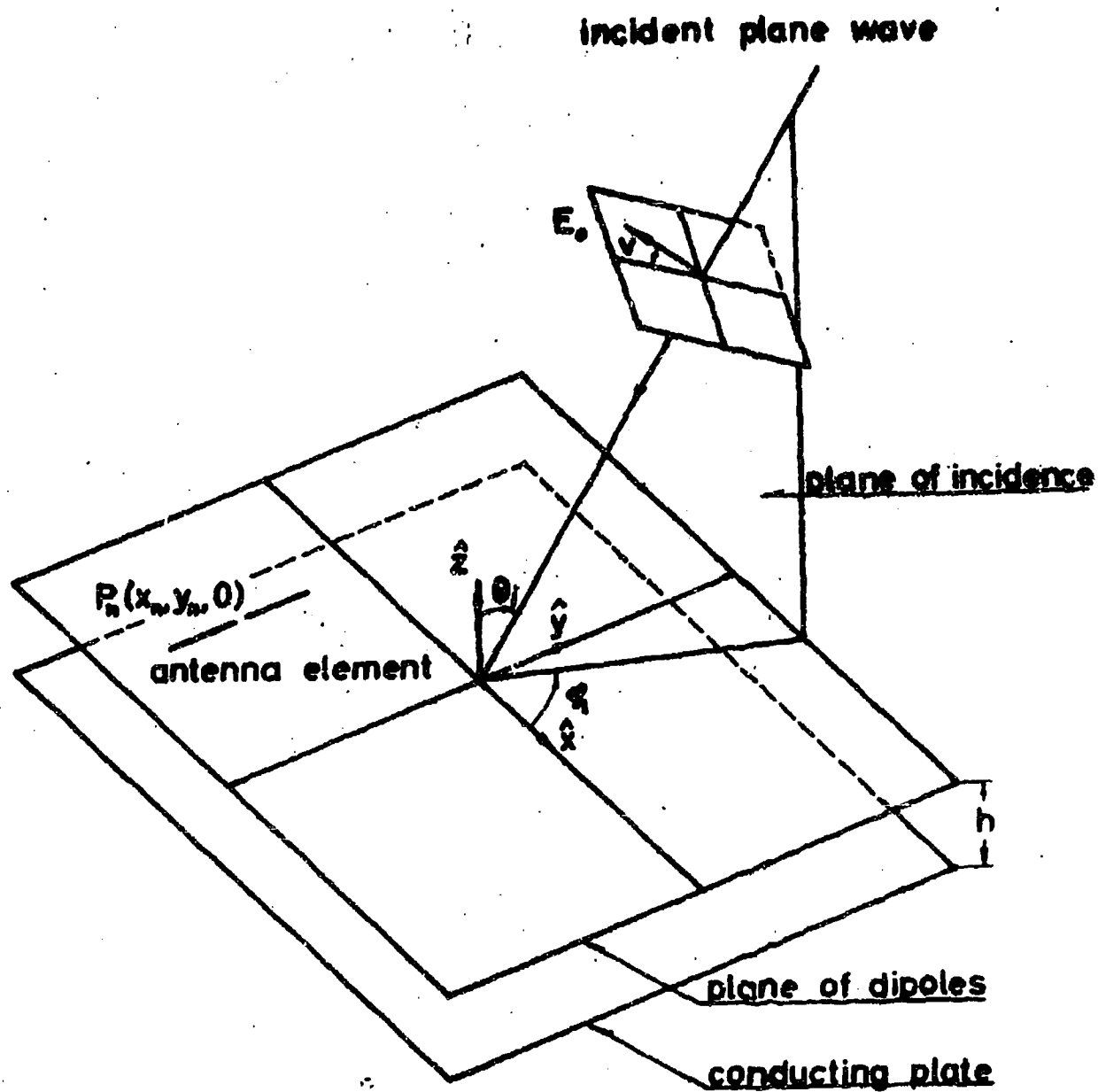


Fig 4. Square Van Atta reflector with conducting plate.

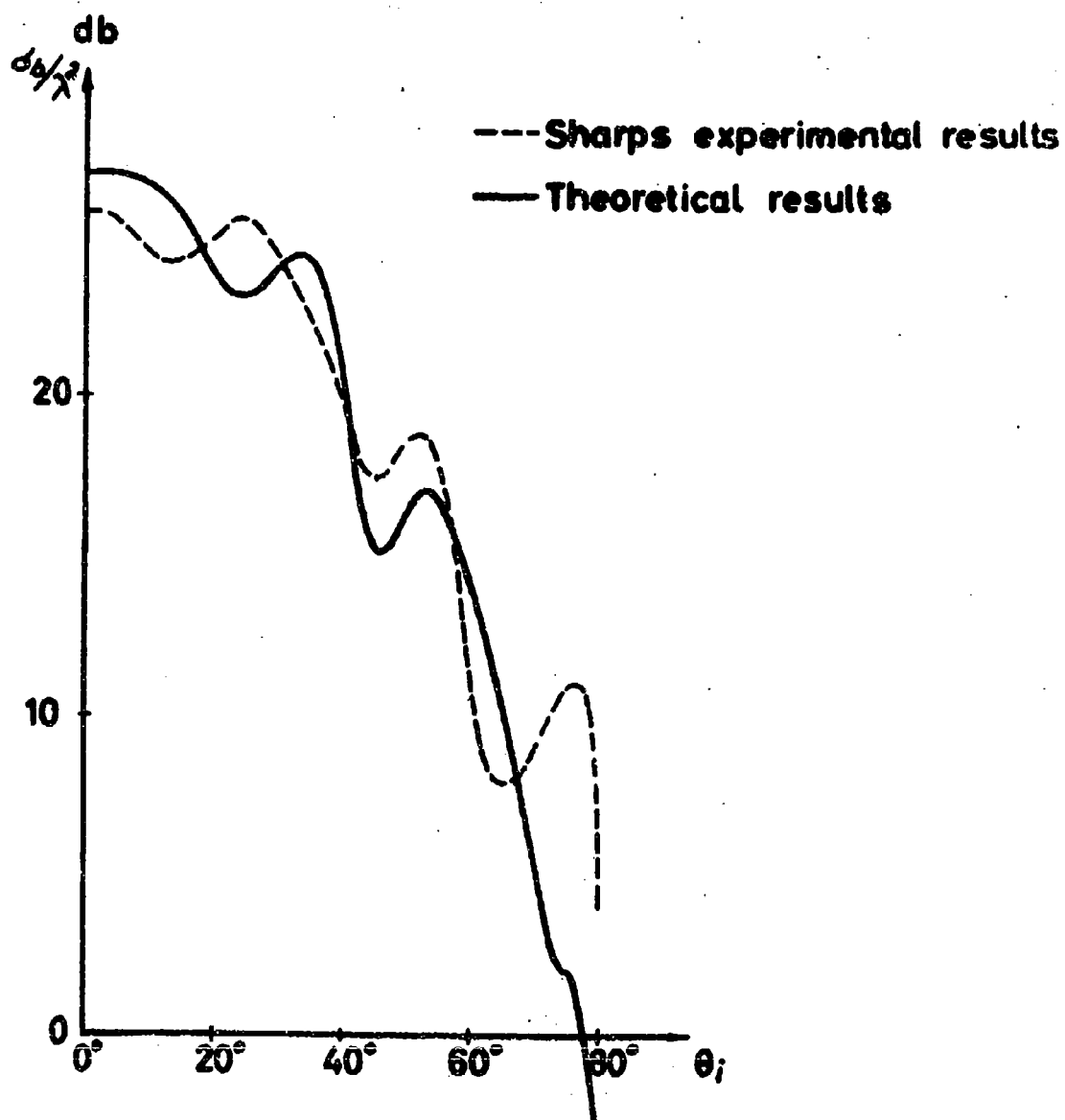


Fig.5. Normalized back-scattering cross section of 16 element square Van Atta reflector with conducting plate.
 $a=0.41\lambda$, $d=0.6\lambda$, $h=0.25\lambda$, $Z_0=73\text{ohms}$,
 $\phi_i=0^\circ$, $\psi=90^\circ$.

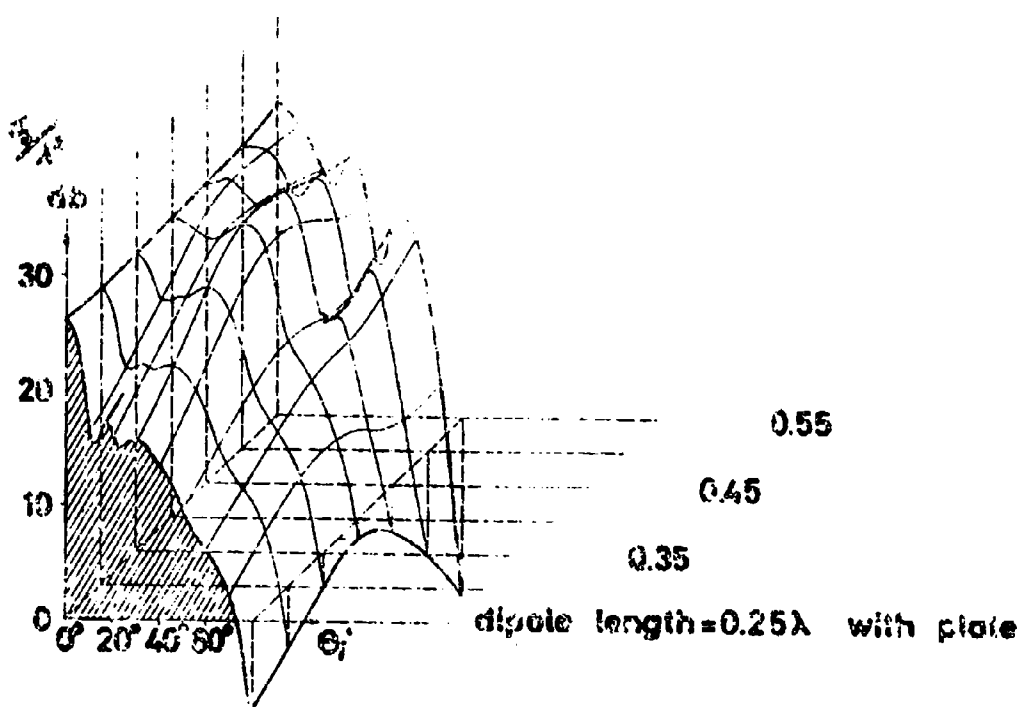
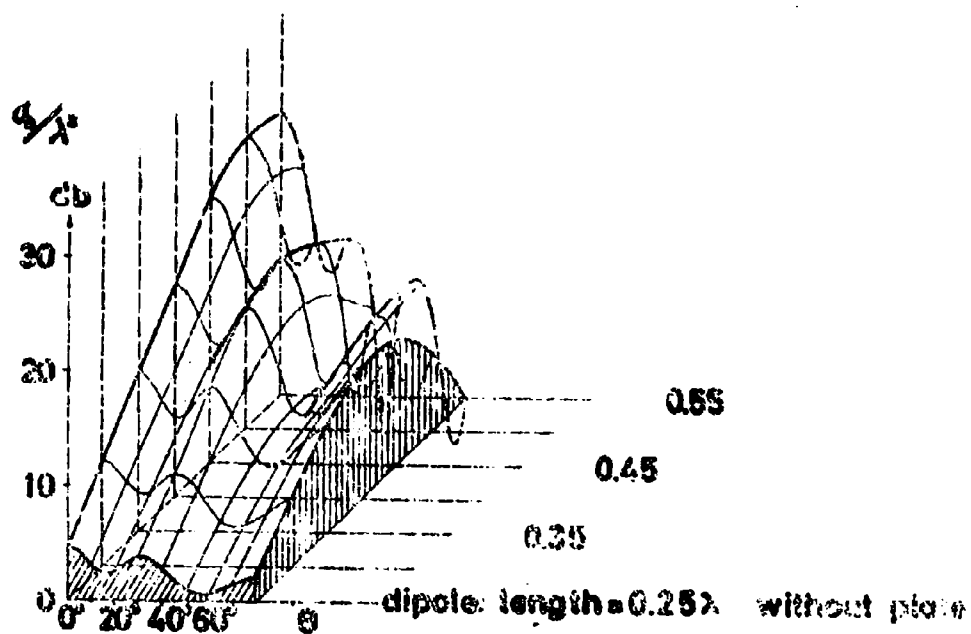
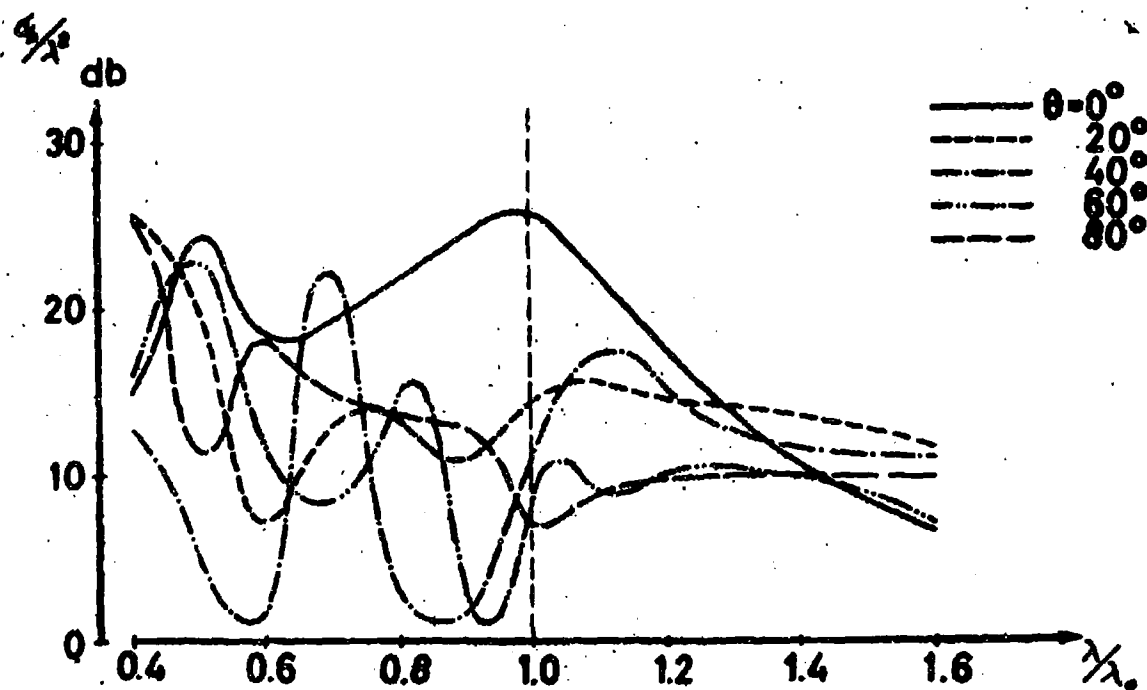
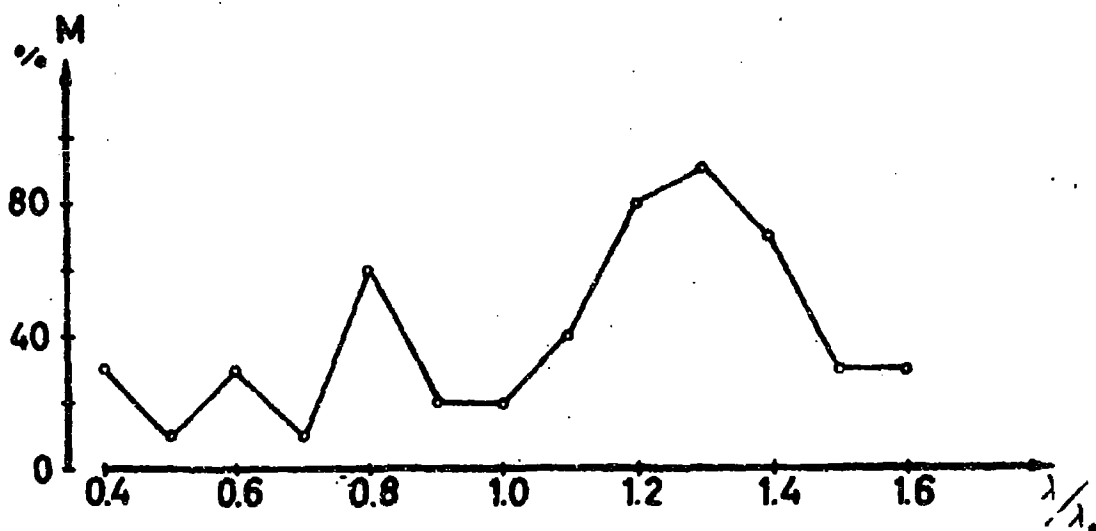


Fig.6. Normalized back-scattering cross section
as a function of the length of the dipoles.
 $N=16$ elements, $a=0.41\lambda$, $d=0.6\lambda$, $h=0.25\lambda$
 $Z_0=73$ ohms, dipole radius $=0.015\lambda$



a. Bandwith properties.



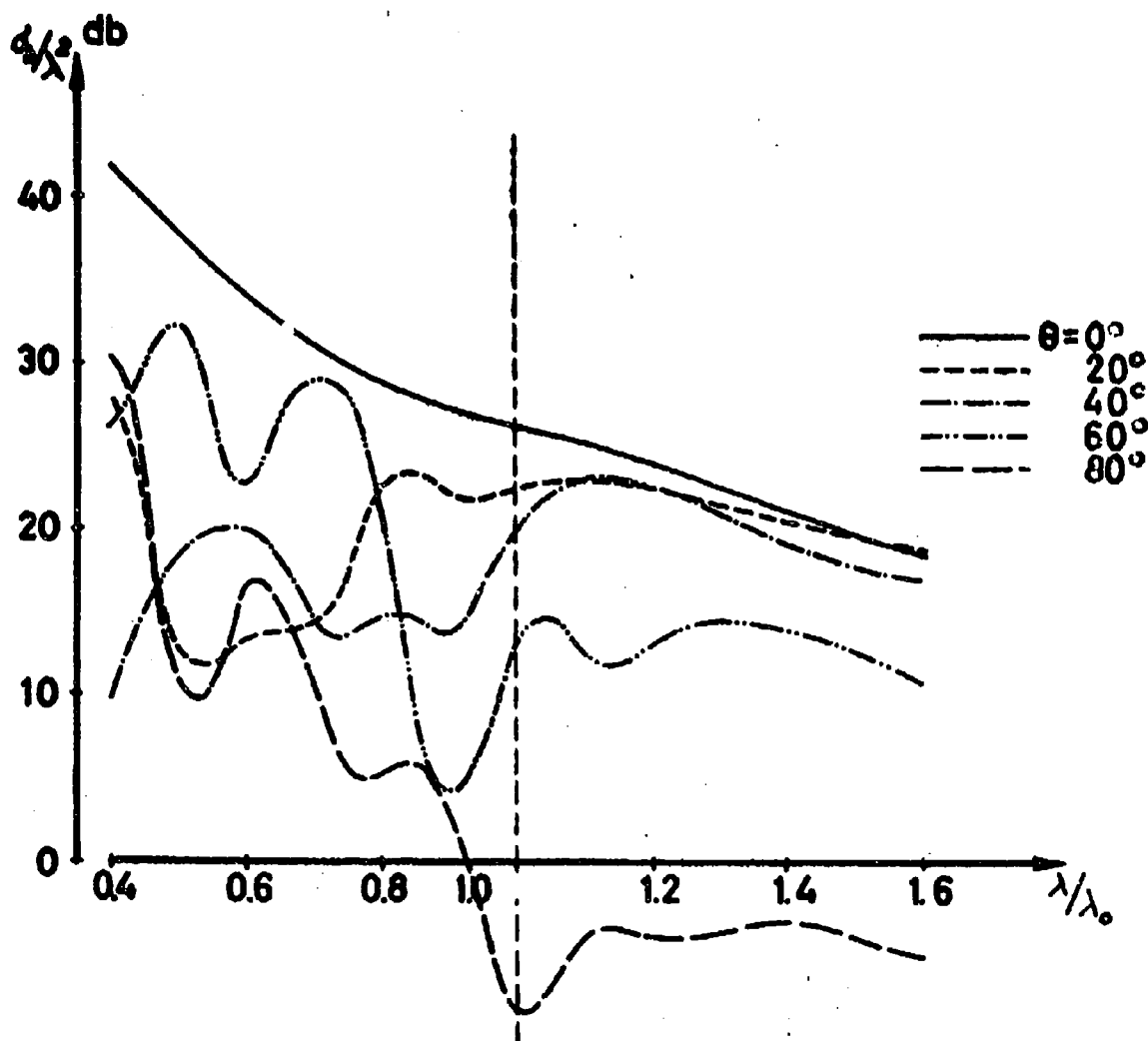
b. Measure M of the retrodirective effect.

Fig.7.16 element square Van Atta reflector without plate.

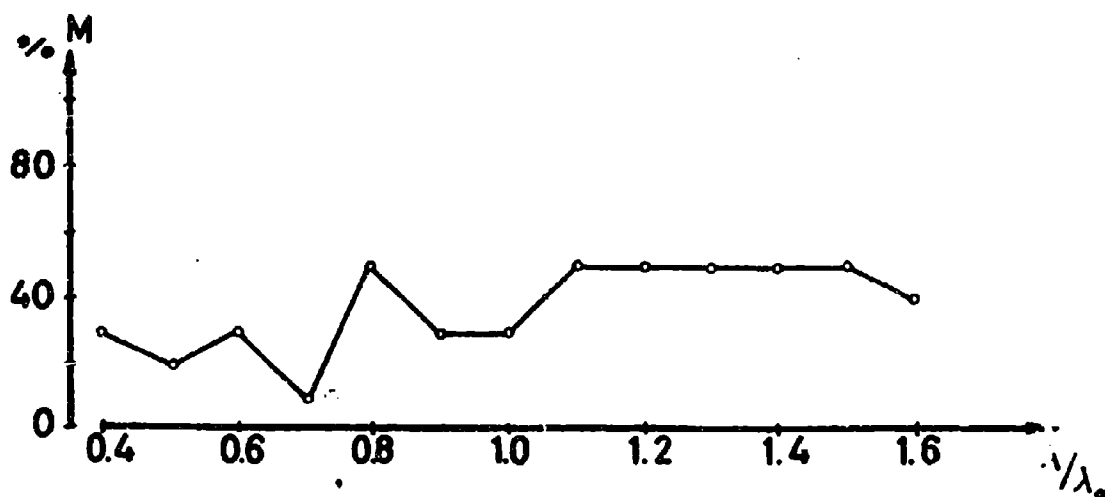
$a = 0.41\lambda_0$, $d = 0.6\lambda_0$, $Z_0 = 73\text{ohms}$

dipole length $= 0.5\lambda_0$, dipole radius $= 0.015\lambda_0$

$\phi = 0^\circ$, $\nu = 90^\circ$



a. Bandwidth properties.



b. Measure M of the retrodirective effect.

Fig.8.16 Element square Van Atta reflector with plate.

$a = 0.41\lambda_0$, $d = 0.6\lambda_0$, $h = 0.25\lambda_0$, $Z_0 = 73$ ohms

dipole length $= 0.5\lambda_0$, dipole radius $= 0.015\lambda_0$

$\varphi_i = 0^\circ$, $\nu = 90^\circ$

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13 ABSTRACT The Van Atta reflector was first described in a patent by Dr. L.C. Van Atta in 1959. The advantage of this passive reflector type should be that the reradiated field has a maximum back in the direction of arrival of the primary plane wave. Since this retrodirective effect of the reflector might be of great importance if used as a navigational aid in the air or at sea, it seemed worth while to carry out a theoretical investigation of such reflectors, especially since only experimental investigations had been made before this contract was initiated. The work performed under the contract dealt mainly with theoretical and numerical investigations of Van Atta reflectors consisting of dipoles. A survey of the literature concerning active or passive Van Atta reflectors has been made. Both a linear and a two-dimensional plane Van Atta reflector has been investigated numerically and a theory for arbitrary Van Atta reflectors has been developed. An experimental investigation of a linear Van Atta reflector was carried out and the results compared with the theoretical results.		

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KEY WORDS	LINK A		LINK B		LINK C	
	ROLE	WT	ROLE	WT	ROLE	WT
Van Atta reflector Reflector Array Adaptive Antenna System Dipoles Mutual Impedance Image Theory Conducting Plate Bandwidth	<div style="font-size: 4em; opacity: 0.5;">DELETED</div>					

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